



GCR Anisotropy Effects on Dose Measurements with MTR/DOSTEL

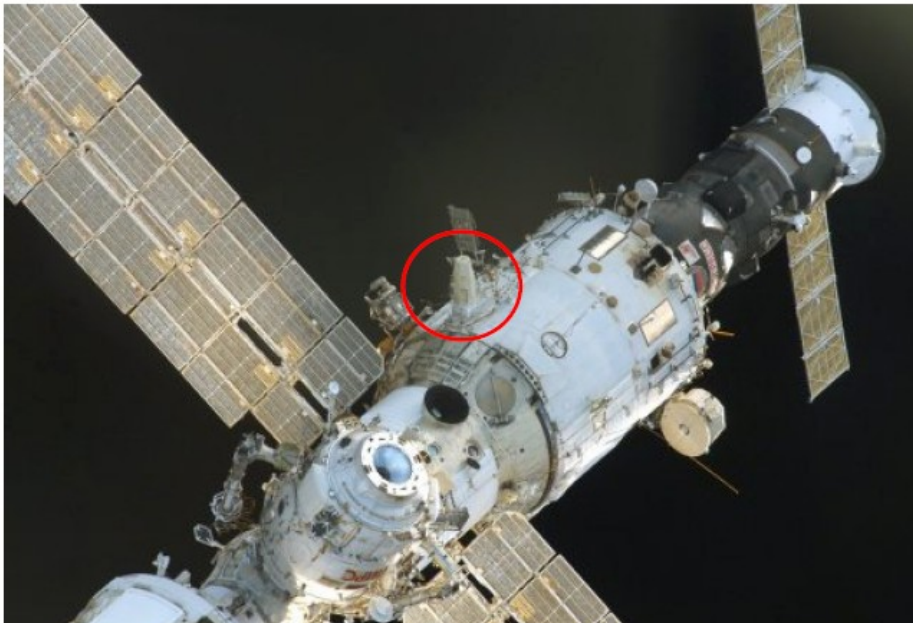
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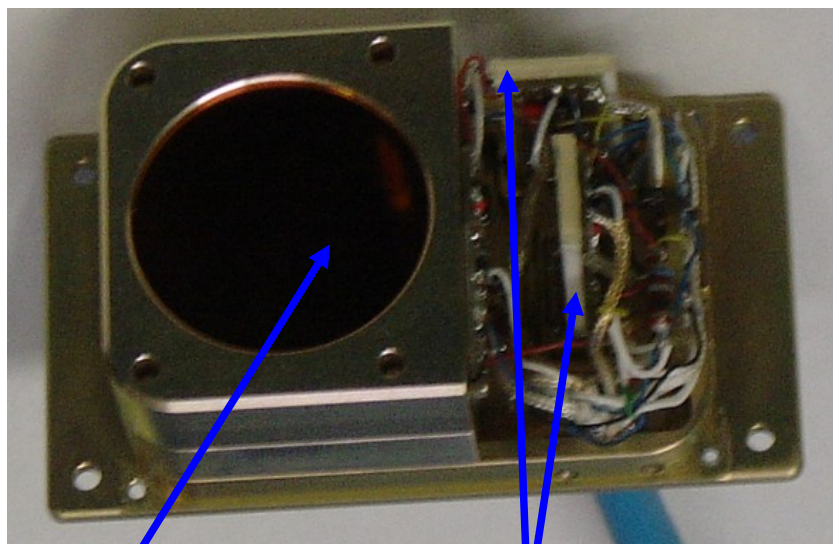
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- Introduction
- Motivation
- Comparison of DOSTEL count rates with Geometric factors (Sullivan)
- Geometric Monte Carlo Model
- Differences between single and telescope detector for different radiation fields
- Summary / Conclusion

- MTR 1**: Feb 2004 – Aug. 2005
outside the ISS
- MTR 2B: Oct. 2007-Nov. 2008
inside the ISS

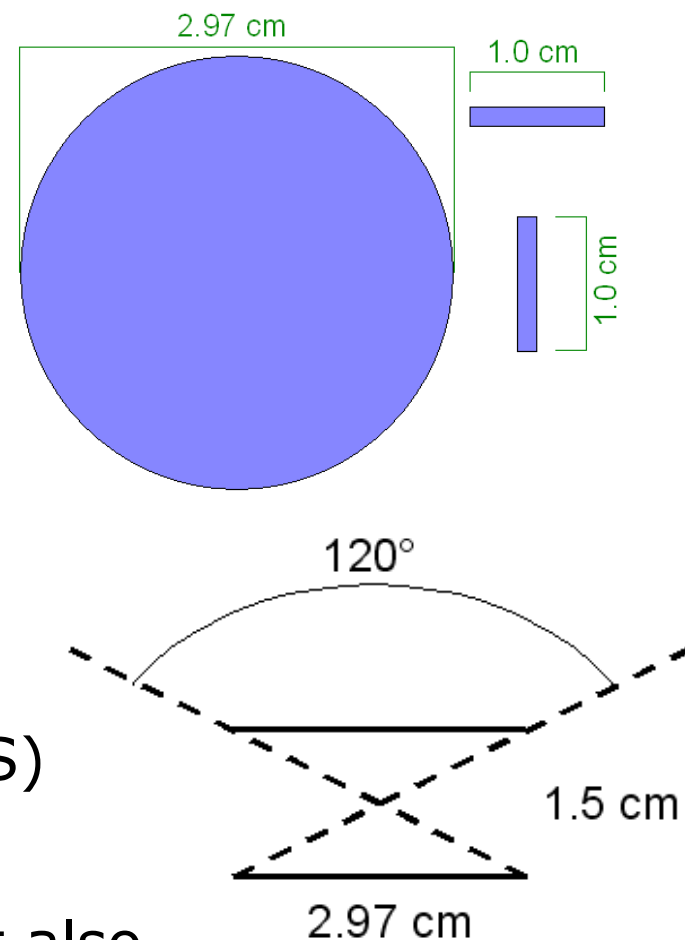




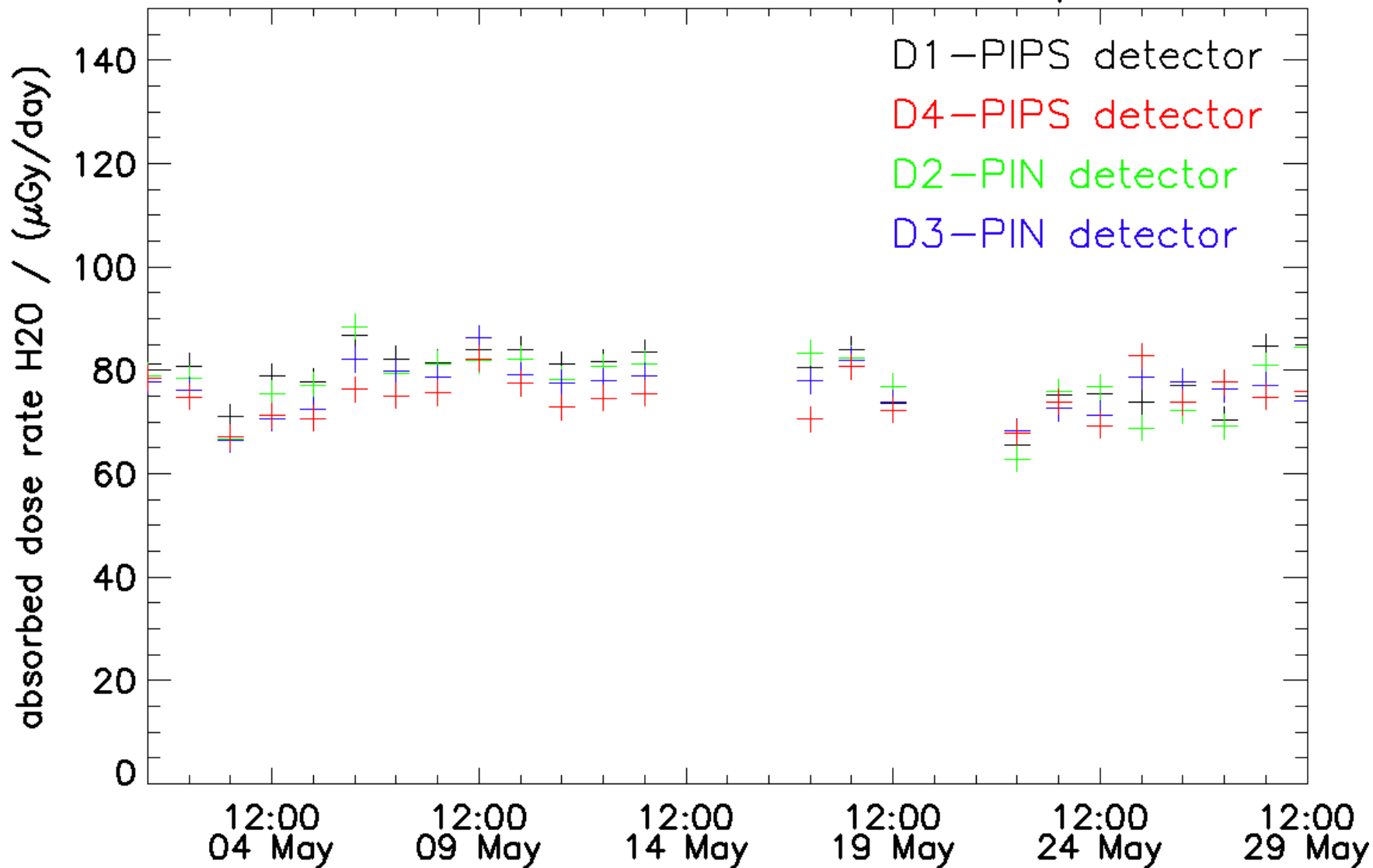
Canberra PIPS
behind capton foil

HAMAMATSU
PIN diodes

- Two silicon detectors (Canberra PIPS) in a telescope geometry (D1,D4).
- PIPS operate in telescope mode, but also as single detectors. (6 hour change)
- In addition 2 Hamamatsu PIN diodes (D2,D3) perpendicular to the PIPS.



MTR 1 abs. dose rate GCR comp.



- The DOSTEL changes every 6 hours between single and telescope mode.

Single mode: PIPS and PINS act as single detectors

Telescope mode: PIPS act as telescope and PINS as single Detectors

- This allows a comparison of single and telescope detector measurements of the GCR component with DOSTEL.
- For this, measured count and dose rate ratios, obtained from the energy deposition spectra, were compared to theoretical ratios between single and telescope detectors.

1. *Circular symmetry*: For a telescope with two circular detectors of radii R_1 and R_2 respectively

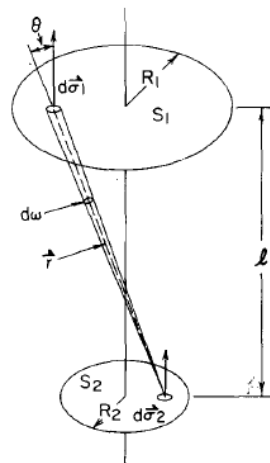


Fig. 2. An ideal cylindrically symmetric telescope with two circular detectors.

(cf. fig. 2), the geometrical factor can be evaluated by direct integration of eq. (5). Whence,

$$G = \frac{1}{2} \pi^2 [R_1^2 + R_2^2 + l^2 - \{(R_1^2 + R_2^2 + l^2)^2 - 4R_1^2 R_2^2\}^{1/2}] \quad (8)$$

For quick estimation, this exact result can be expanded yielding to the first order:

$$G \geq \frac{A_1 A_2}{R_1^2 + R_2^2 + l^2} \quad (9)$$

It should be noted that eq. (7) holds for all telescopes, whereas eq. (9) is applicable only to two circular-detector telescopes. Further, we can evaluate the

J.D. Sullivan, Geometric factor and directional response of single and multi-element particle telescopes, Nuclear Instruments and Methods, Volume 95, Issue 1, 1 August 1971, Pages 5-11.

PIPS (D1,D4):

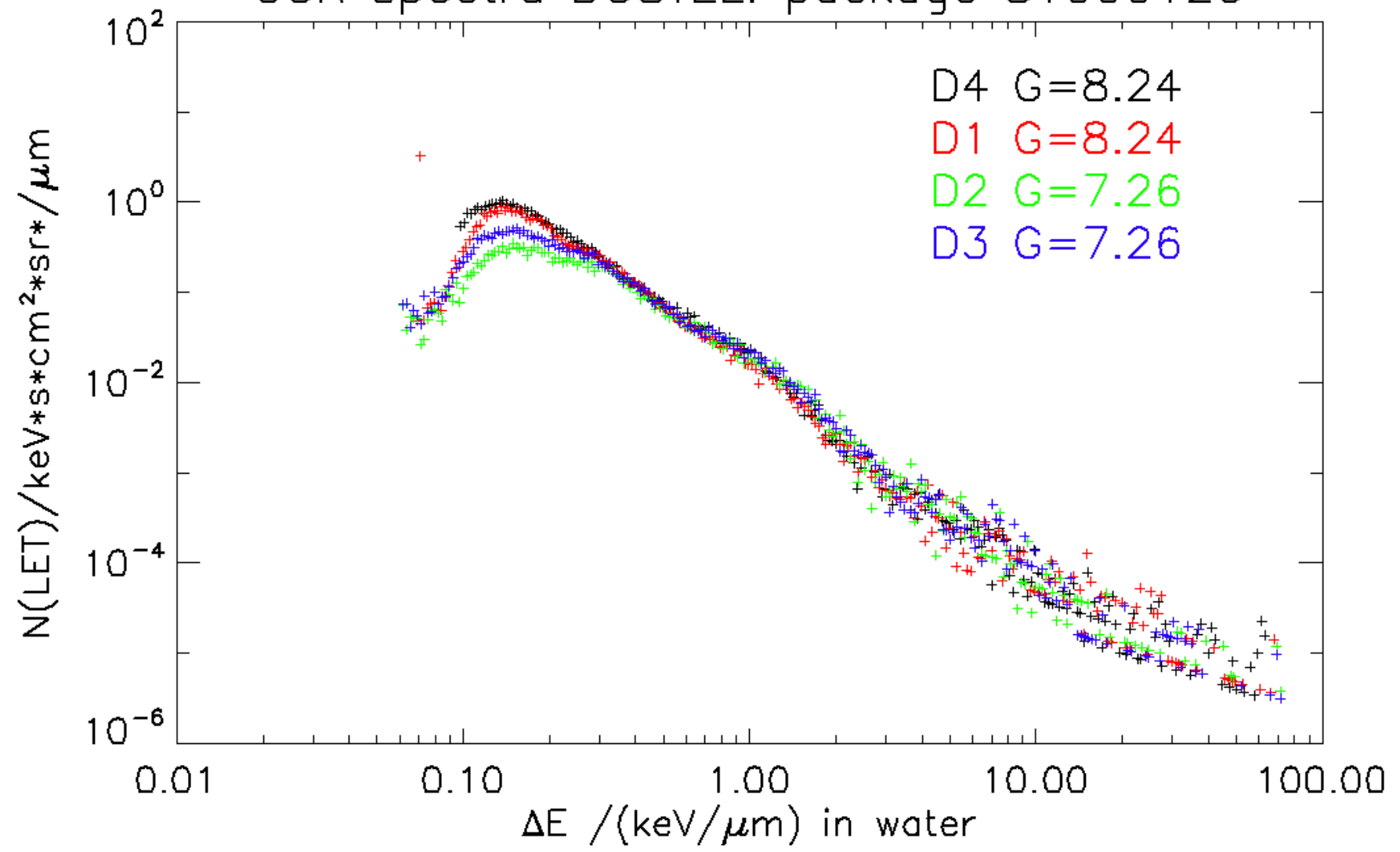
Telescope: $R_1 = R_2 = 1.485$ cm $l = 1.50$ cm $\rightarrow G_{tel} = 8.24$ cm² sr

Single: $R_1 = R_2 = 1.485$ cm $l = 0.0$ cm $\rightarrow G_{sin} = 21.76$ cm² sr

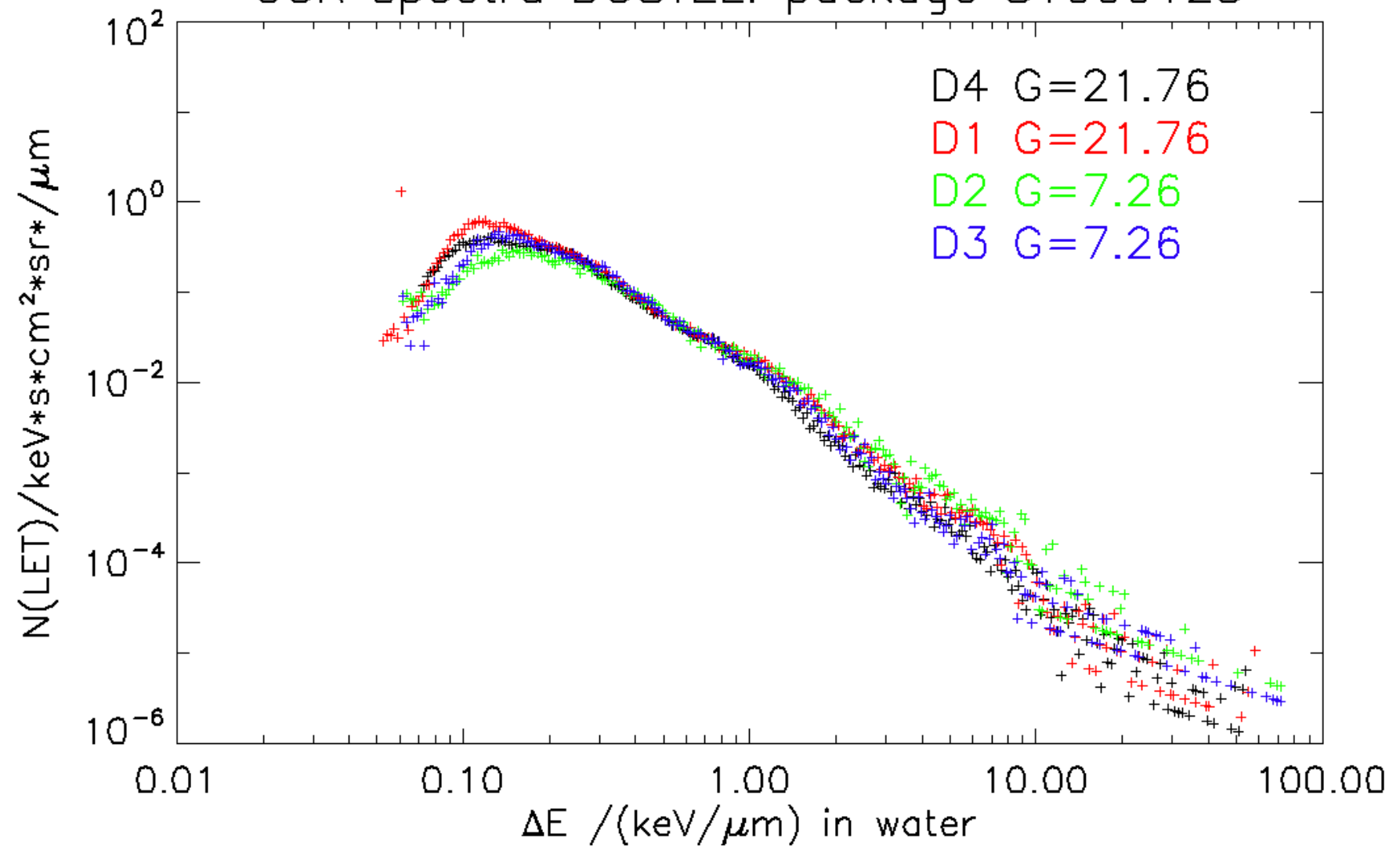
PINS (D2,D3):

$\rightarrow G = 7.26$ cm² sr

GCR spectra DOSTEL: package S1000126



GCR spectra DOSTEL: package S1000125



Upper PIPS detector:

$$D1: \frac{\text{Count rate single mode spectra}}{\text{Count rate telescope mode spectra}} = \underline{2.13}$$

Lower PIPS detector:

$$D4: \frac{\text{Count rate single mode spectra}}{\text{Count rate telescope mode spectra}} = \underline{1.91}$$

PIN Diodes:

$$D2: \frac{\text{Count rate single mode spectra}}{\text{Count rate telescope mode spectra}} = 0.966$$

$$D3: \frac{\text{Count rate single mode spectra}}{\text{Count rate telescope mode spectra}} = 1.00$$

All spectra from April to June 2004 were used, to avoid geo-magnetic differences for the different spectra types.

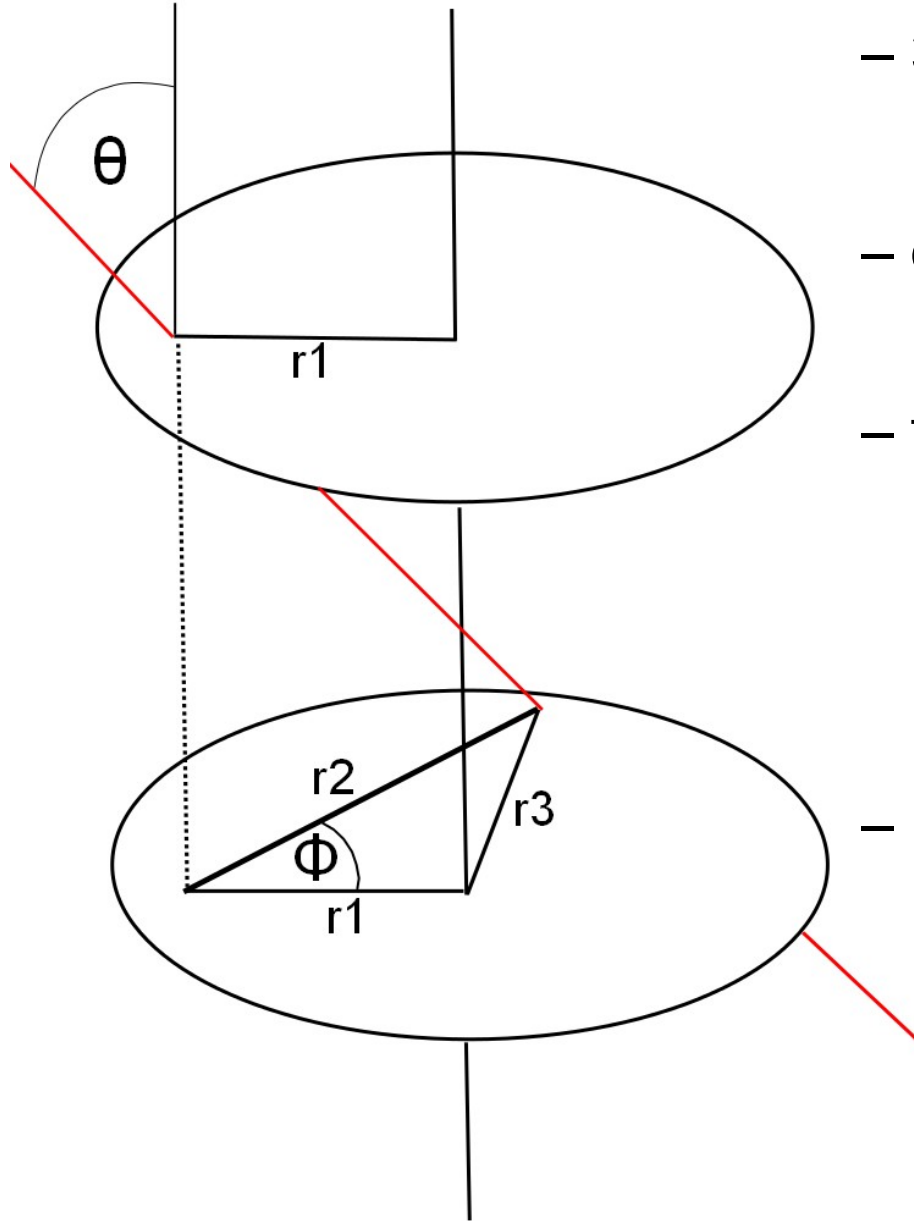
- Theoretical count rate ratio for an isotropic field between a single and a telescope detector (Sullivan):

$$\text{Ratio} = 21.76 \text{ cm}^2 \text{ sr} / 8.24 \text{ cm}^2 \text{ sr} = \underline{2.64}$$

- Measured count rate ratio between single and telescope mode for the lower PIPS detector (D4):

$$\text{Ratio} = \underline{1.91}$$

- To explain these differences, a simple Monte-Carlo-Model was used to obtain the theoretical count rate ratios between a single and a telescope detector for different anisotropic GCR distributions.



– 3 independent series of random numbers between 0 and 1 (a,b,c)

– distance to center of first detector
 $r1 = \sqrt{a} * R$, $R = \text{Radius of detector}$

– theta angle:

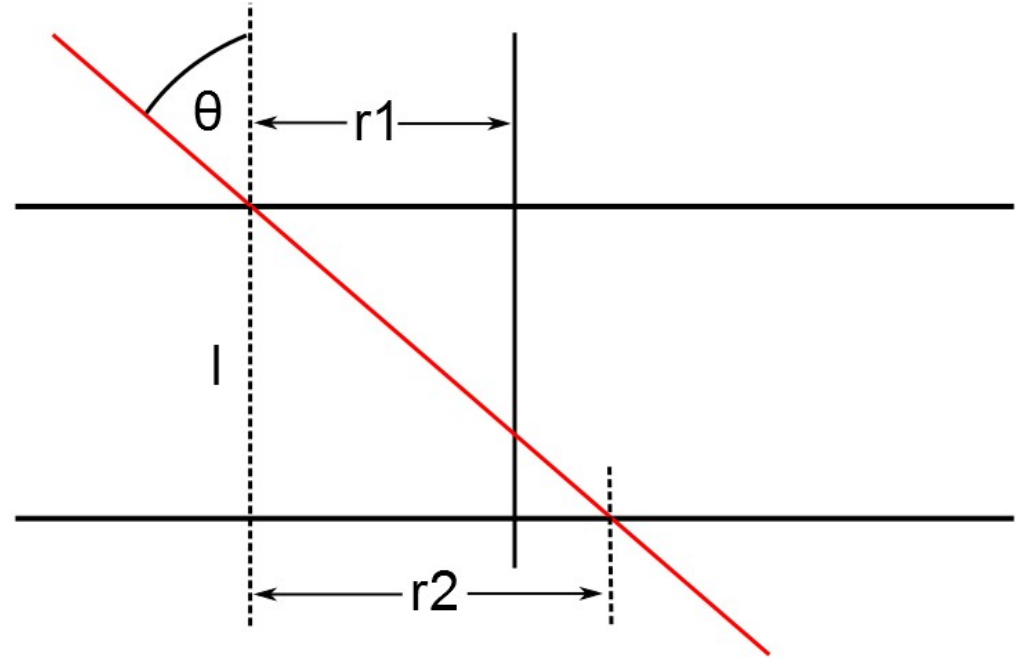
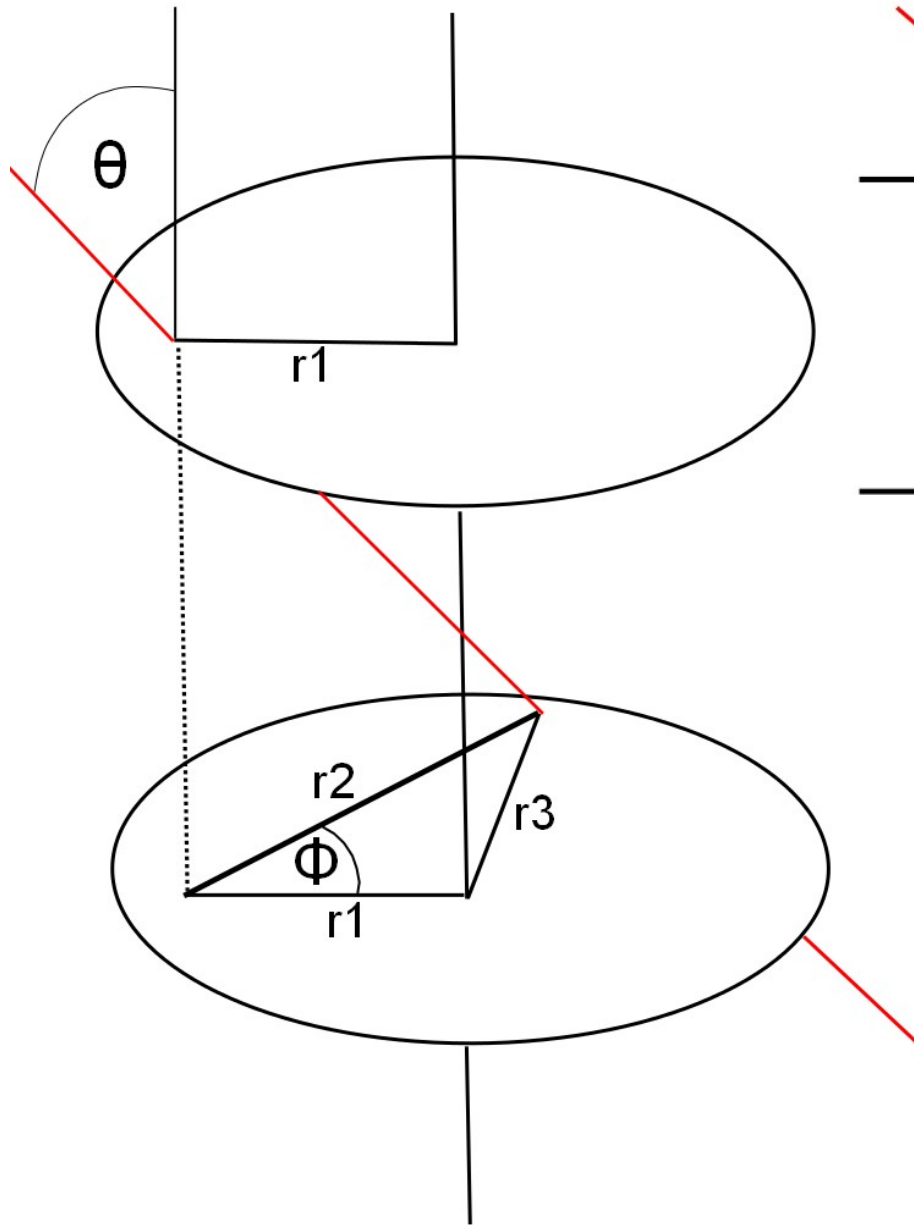
$$P = \int \cos(\theta) * \sin(\theta) d\theta = 0.5 \cos^2(\theta)$$

$$\theta \approx \arccos(\sqrt{P})$$

$$\theta = \arccos(\sqrt{b}), (\text{isotropic})$$

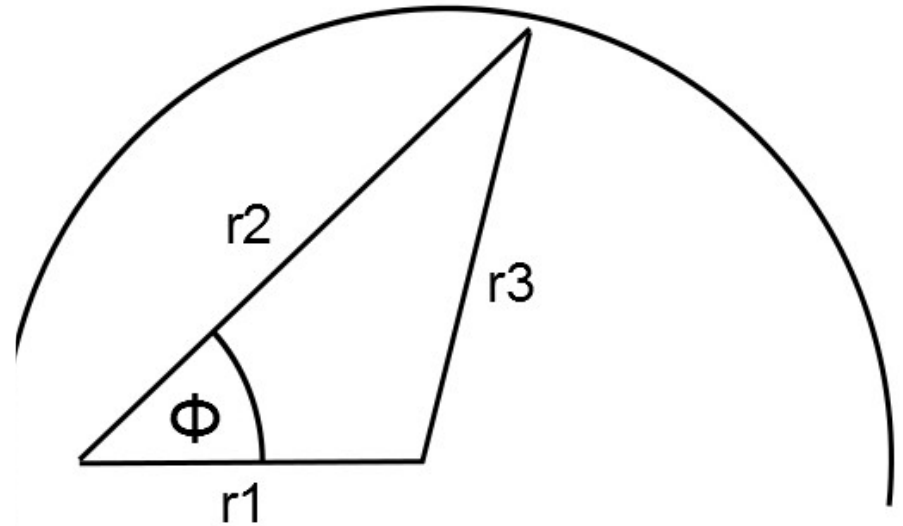
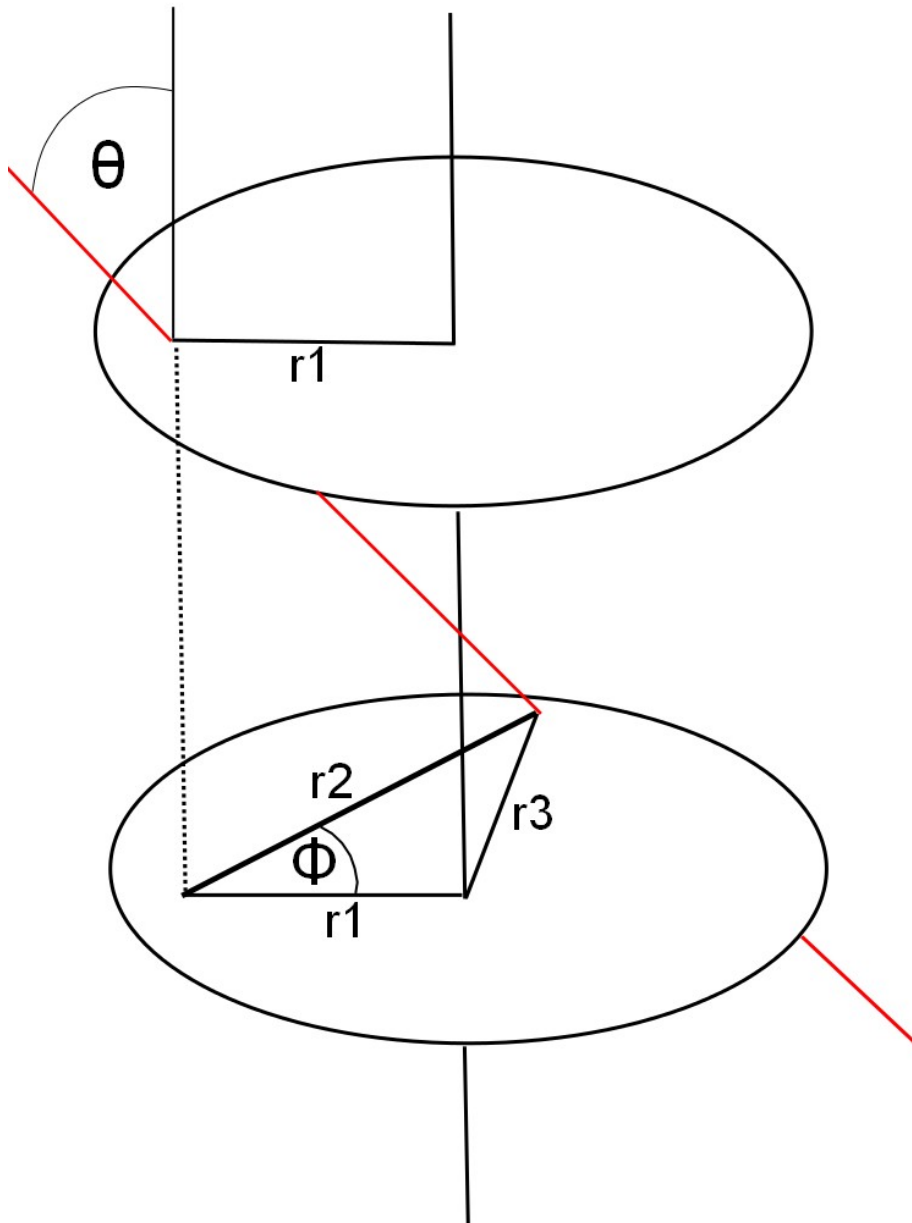
– phi angle:

$$\Phi = 2 * \Pi * c$$



– r_2 : projected distance between points of impact on first and second detector

$$r_2 = l * \tan(\theta)$$



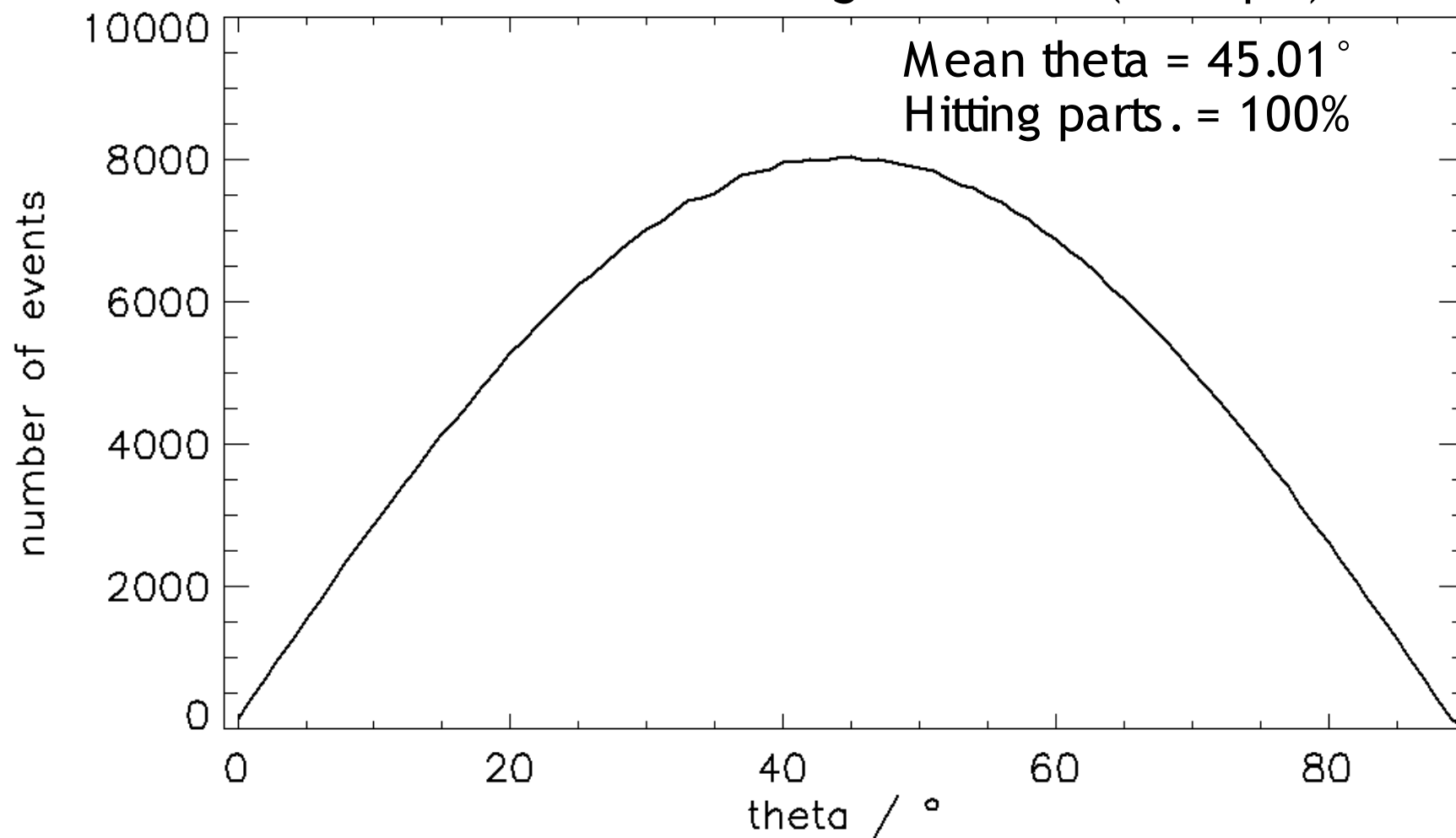
– r_3 : distance between point on- and center of- second detector

– with law of cosine:

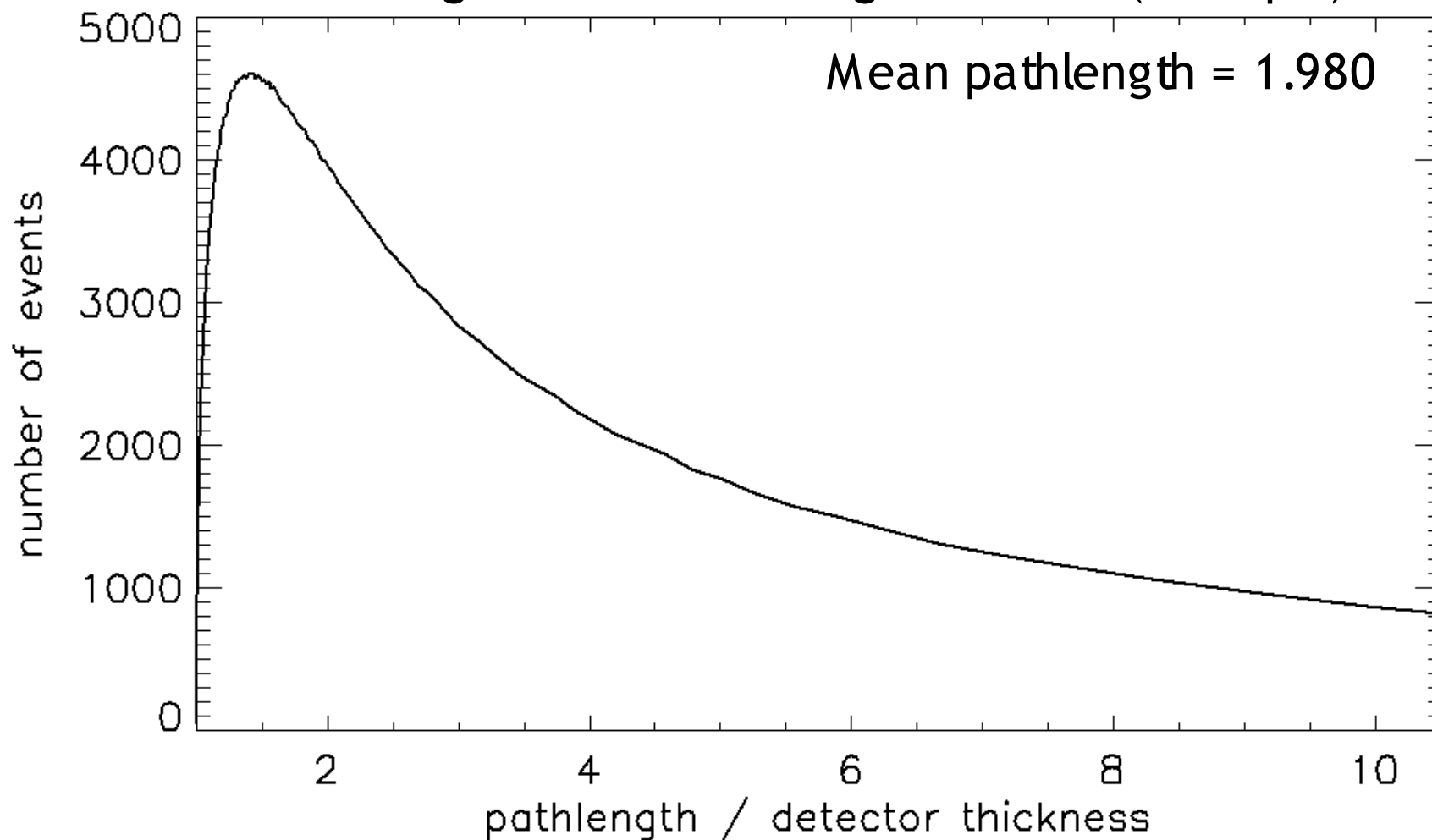
$$r_3 = \sqrt{r_1^2 + r_2^2 - (2 * r_1 * r_2 * \cos(\Phi))}$$

- If r_3 is less than R (1.485 cm), then the particle hits both detectors and only these particles are counted.
- For telescope detector a distance of $l=1.50$ cm between the detectors was used.
- For single detector mode the distance was set to 0.0 cm
- The number of particles hitting the detector are counted, to get the ratio of the geometric factors between single and telescope detector.
- The mean theta angles and the corresponding mean path length in the detectors were analyzed.

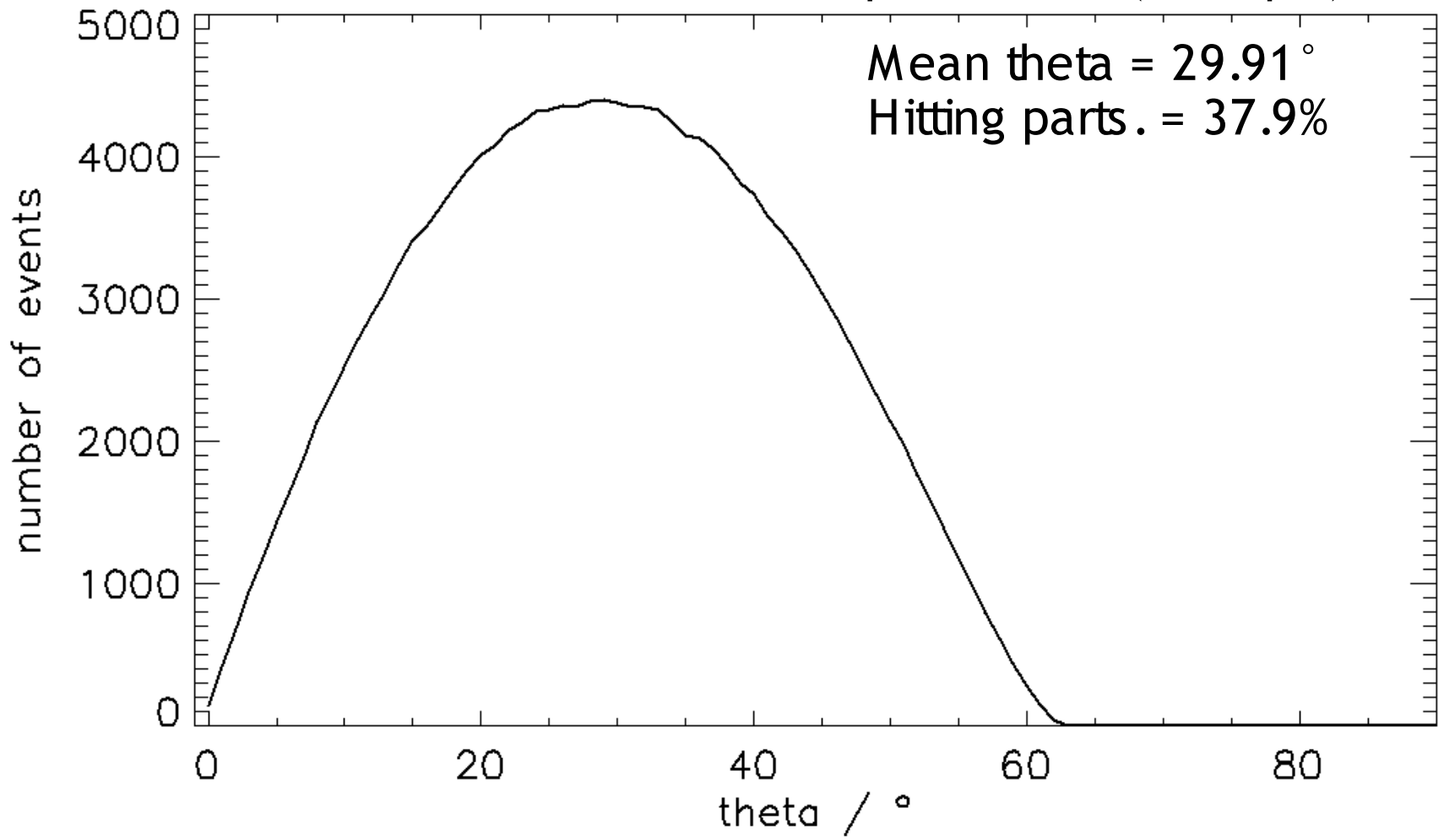
Theta in DOSTEL single detector (isotropic)



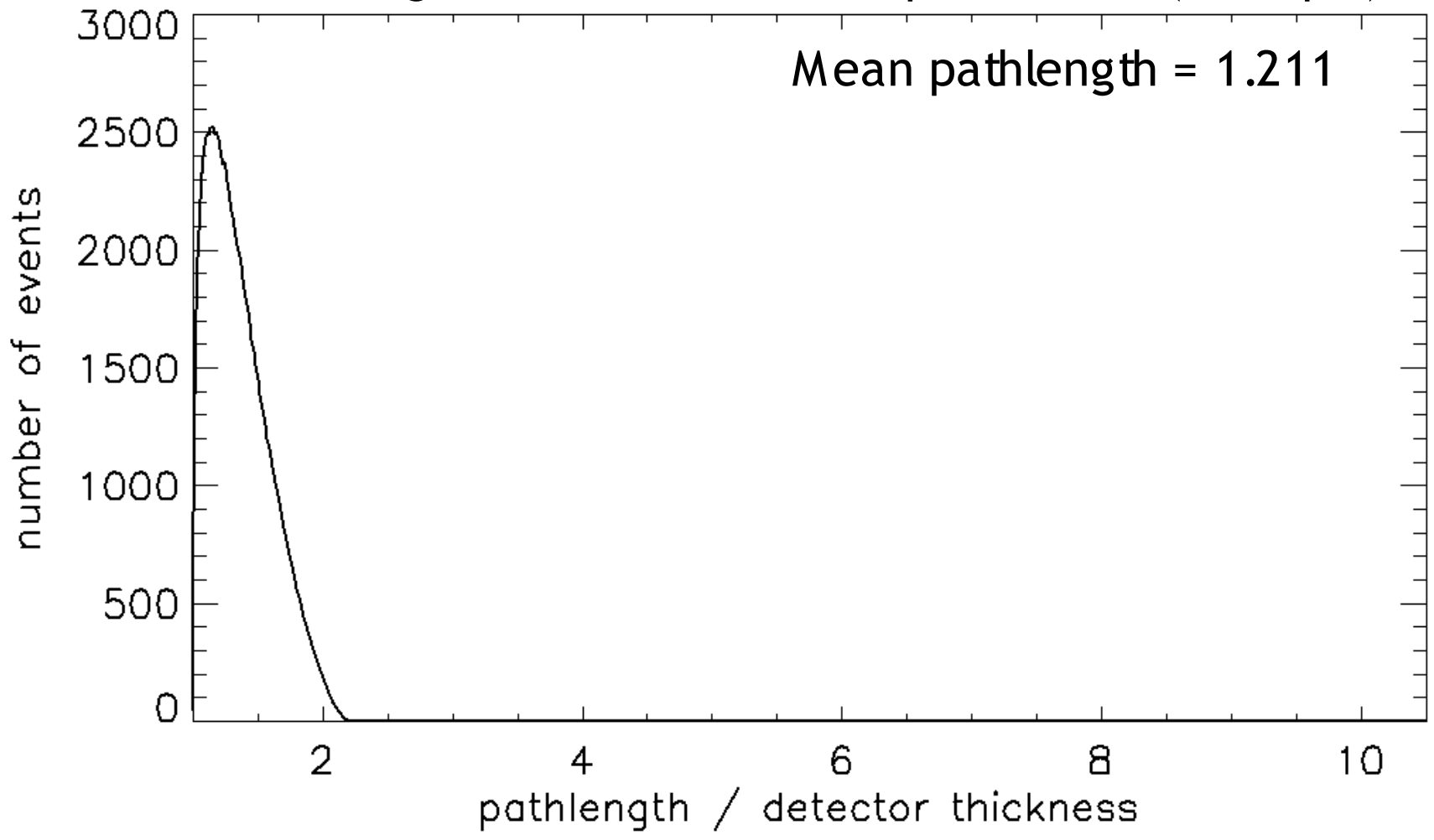
Pathlength in DOS TEL single detector (isotropic)



Theta in DOS TEL telescope detector (isotropic)



Pathlength in DOSTEL telescope detector (isotropic)



- Geometric factors calculated with Sullivan formula:

$$\text{Telescope mode: } G_{\text{Tel}} = 8.24 \text{ cm}^2 \text{ sr}$$

$$\text{Single mode: } G_{\text{Sin}} = 21.76 \text{ cm}^2 \text{ sr}$$

$$G_{\text{Sin}} / G_{\text{Tel}} = \underline{2.64}$$

- Geometric factor calculated with model for isotropic field:

$$\text{hits}_{\text{Sin}} / \text{hits}_{\text{Tel}} = 100.0\% / 37.88\% = \underline{2.64}$$

- Count rate ratio DOSTEL:

$$N_{\text{sin}} / N_{\text{tel}} = \underline{1.91}$$

Isotropic (cos¹):

$$P = \int \cos(\theta) * \sin(\theta) d\theta = 0.5 \cos^2(\theta)$$

$$\theta \approx \arccos(\sqrt{P})$$

$$\theta = \arccos(\sqrt{b})$$

Cos²:

$$P = \int \cos^2(\theta) * \sin(\theta) d\theta = \frac{1}{3} \cos^3(\theta)$$

$$\theta \approx \arccos(\sqrt[3]{P})$$

$$\theta = \arccos(\sqrt[3]{b})$$

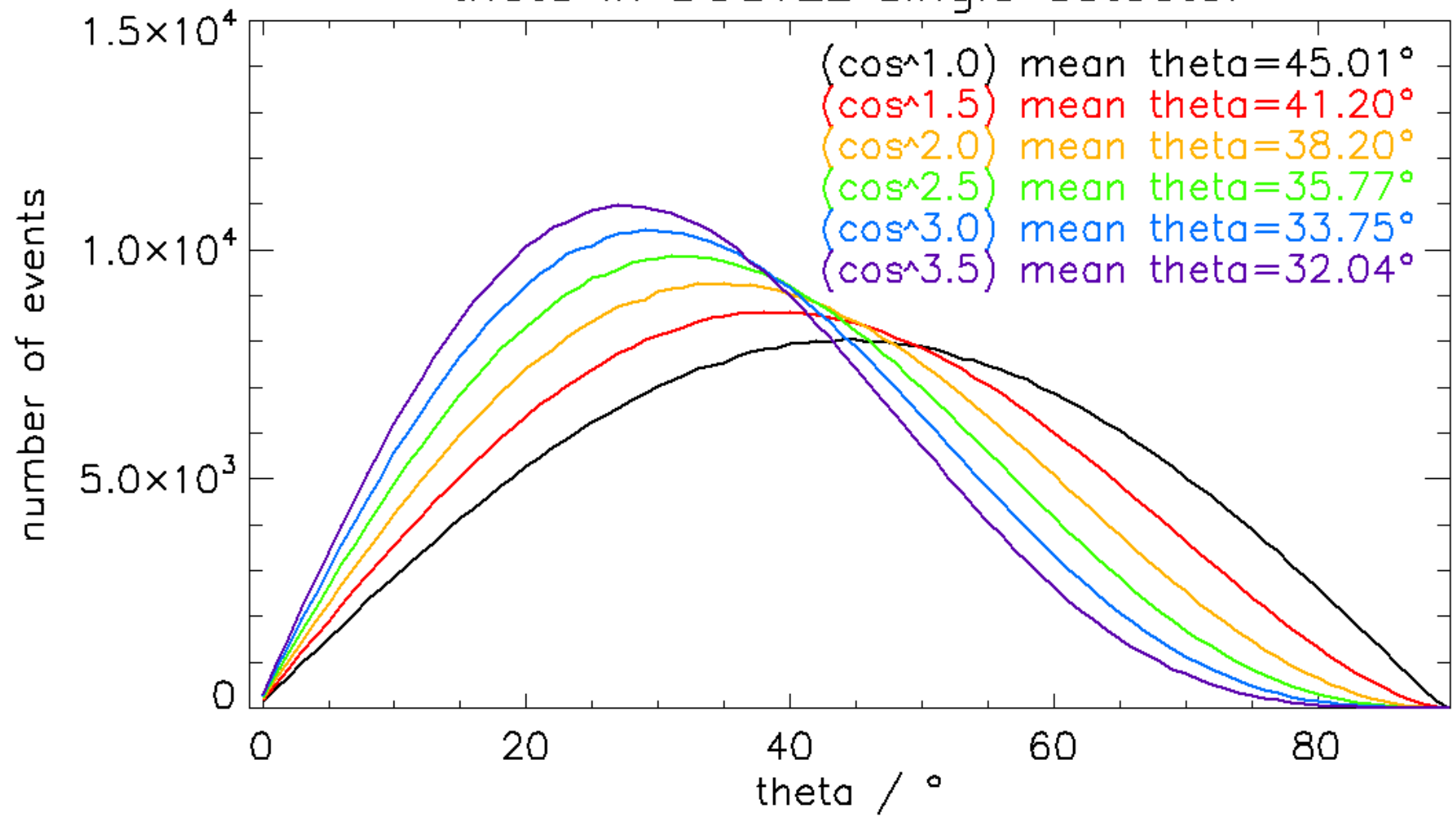
Cosⁿ:

$$P = \int \cos^n(\theta) * \sin(\theta) d\theta = \frac{1}{n+1} \cos^{n+1}(\theta)$$

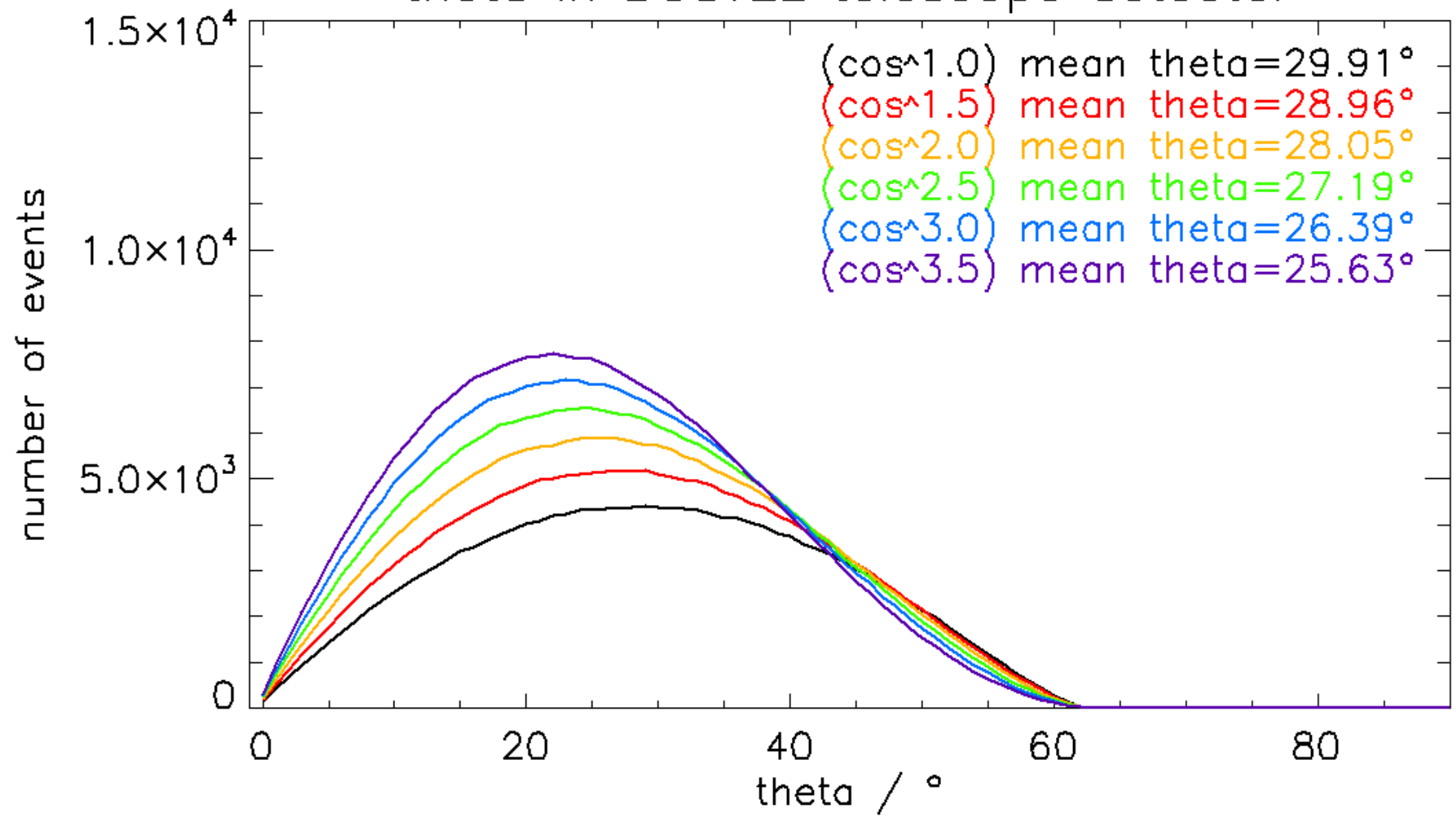
$$\theta \approx \arccos(\sqrt[n+1]{P})$$

$$\theta = \arccos(\sqrt[n+1]{b})$$

theta in DOSTEL single detector



theta in DOSTEL telescope detector



Simulated ratio of particle numbers in single and telescope detectors:

$$\text{Cos}^1: \quad 100\% / 37.9\% = 2.64$$

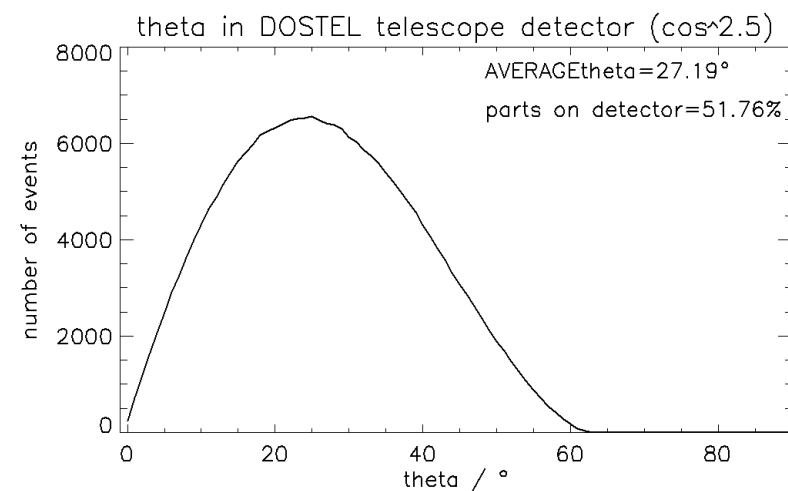
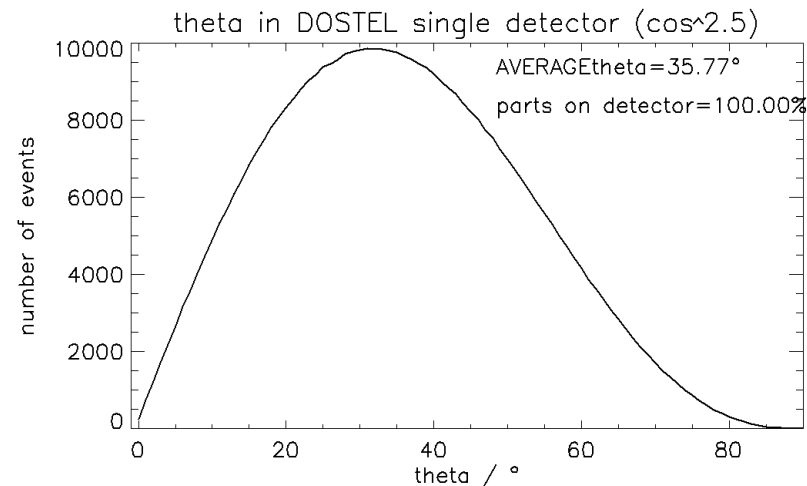
$$\text{Cos}^{1.5}: \quad 100\% / 43.4\% = 2.30$$

$$\text{Cos}^2: \quad 100\% / 48.0\% = 2.08$$

$$\text{Cos}^{2.5}: \quad 100\% / 51.8\% = \underline{1.93}$$

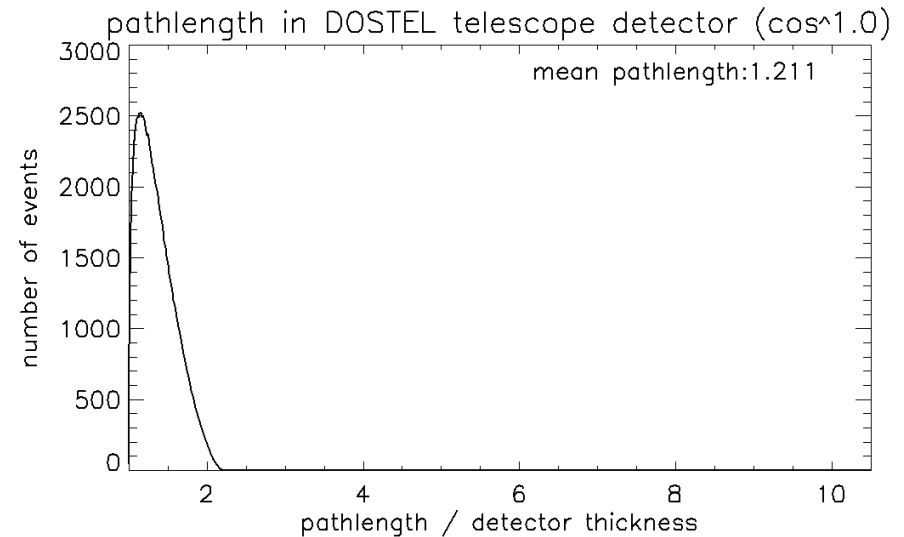
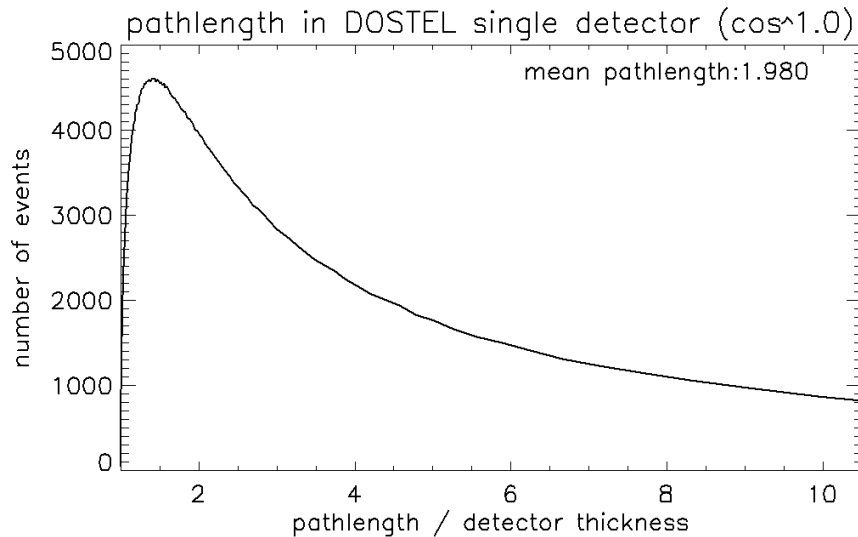
$$\text{Cos}^3: \quad 100\% / 55.0\% = 1.82$$

$$\text{Cos}^{3.5}: \quad 100\% / 57.7\% = 1.73$$



- $\text{Cos}^{2.5}$ is closest to measured ratio (1.91)

- The second measured value to compare with the model is the absorbed dose.
- The absorbed dose is proportional to the number of particles multiplied with the mean particle pathlength.
- The count rate ratio multiplied with the pathlength ratio equals the absorbed dose rate ratio.



Upper PIPS detector:

$$D1: \frac{\text{Absorbed dose in single mode spectra}}{\text{Absorbed dose in telescope mode spectra}} = \underline{2.32}$$

Lower PIPS detector:

$$D4: \frac{\text{Absorbed dose in single mode spectra}}{\text{Absorbed dose in telescope mode spectra}} = \underline{2.35}$$

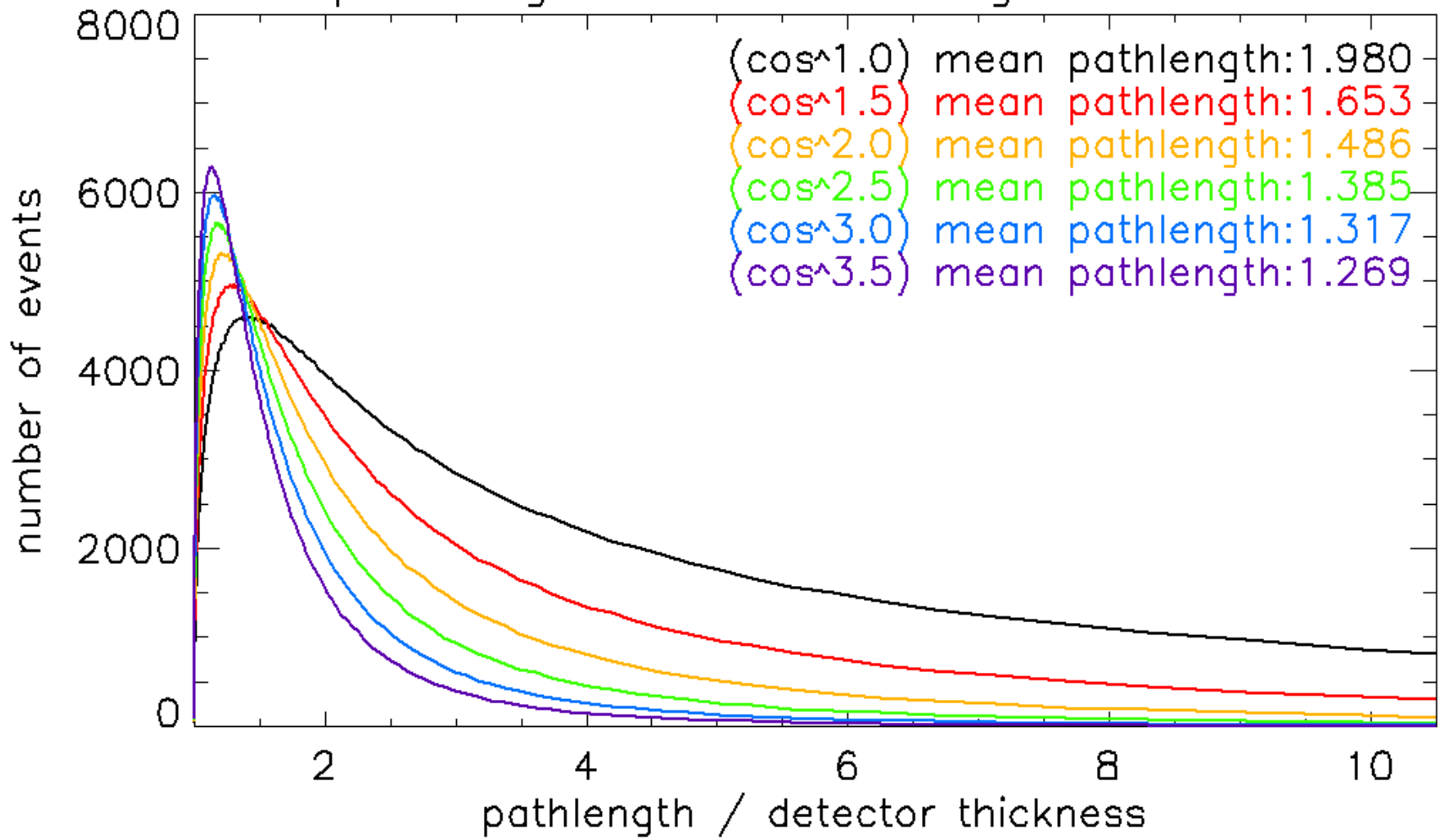
PIN Diodes:

$$D2: \frac{\text{Absorbed dose in single mode spectra}}{\text{Absorbed dose in telescope mode spectra}} = 0.97$$

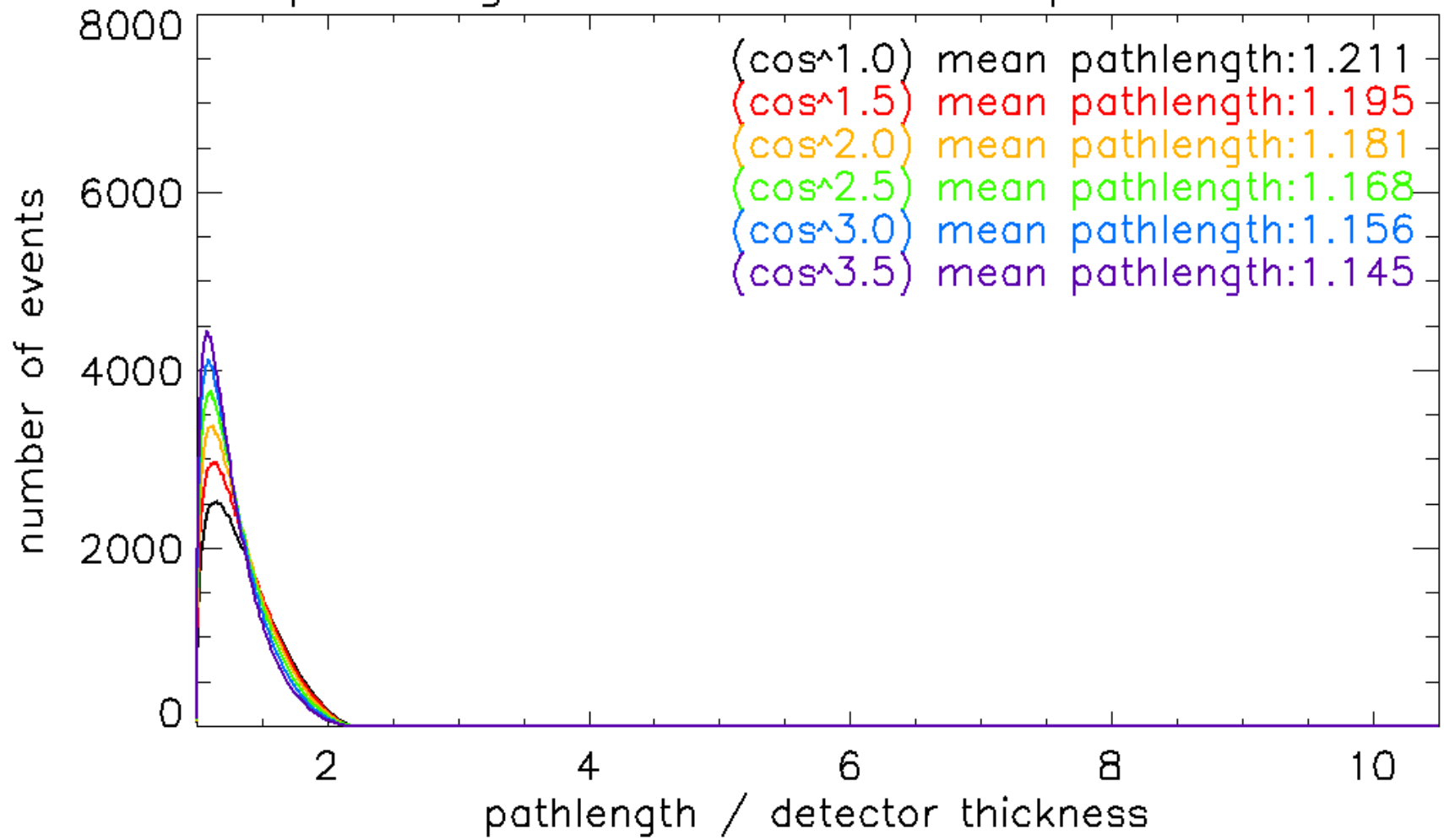
$$D3: \frac{\text{Absorbed dose in single mode spectra}}{\text{Absorbed dose in telescope mode spectra}} = 1.01$$

All spectra from April to July 2004 were used, to avoid geomagnetic differences for the different spectra types.

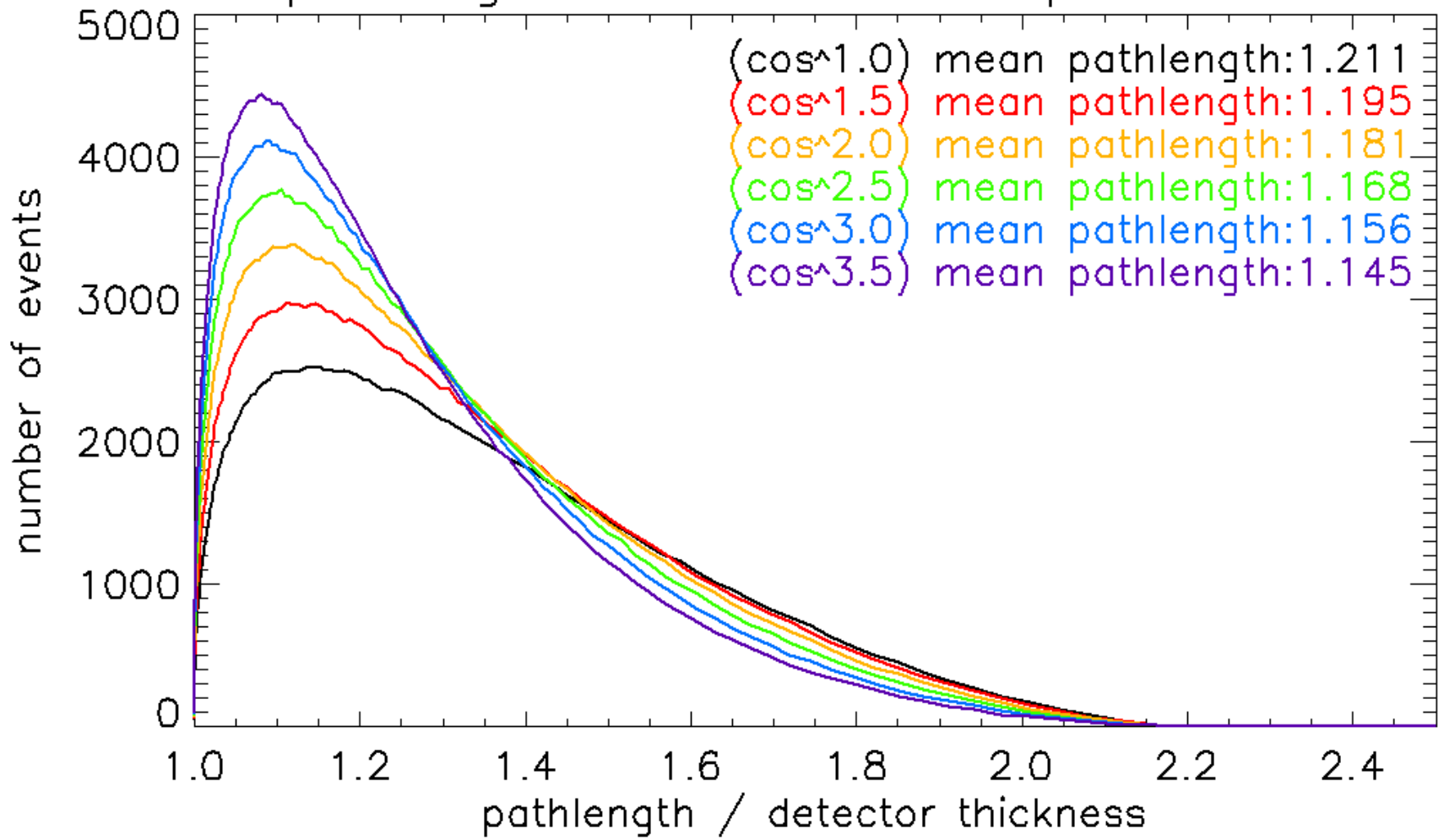
pathlength in DOSTEL single detector



pathlength in DOSTEL telescope detector



pathlength in DOSTEL telescope detector



Simulated ratio of absorbed dose in single and telescope detectors:

$$\text{Cos}^1: \quad 100\% / 37.9\% * 1.98 / 1.21 = 4.32$$

$$\text{Cos}^{1.5}: \quad 100\% / 43.4\% * 1.65 / 1.20 = 3.17$$

$$\text{Cos}^2: \quad 100\% / 48.0\% * 1.49 / 1.18 = 2.63$$

$$\text{Cos}^{2.5}: \quad 100\% / 51.8\% * 1.39 / 1.17 = \underline{2.29}$$

$$\text{Cos}^3: \quad 100\% / 55.0\% * 1.32 / 1.16 = 2.09$$

$$\text{Cos}^{3.5}: \quad 100\% / 57.7\% * 1.27 / 1.15 = 1.91$$

- $\text{Cos}^{2.5}$ is closest to measured ratio (2.32, 2.35) again.

For a radiation field with a $\text{Cos}^2.5$ theta distribution the theoretical ratio of particles in single and telescope mode detectors is:

$$N_{\text{single}} / N_{\text{telescope}} = 100\% / 51.8\% = \underline{1.93}$$

This is in good comparison to the particle number ratio measured with the lower DOS TEL detector (1.91) during MTR1 mission phase.

The corresponding pathlength ratio,

$$\text{Path}_{\text{single}} / \text{path}_{\text{telescope}} = 1.39 / 1.17 = 1.19$$

then leads to a theoretical absorbed dose ratio of:

$$D_{\text{single}} / D_{\text{telescope}} = 1.93 * 1.19 = \underline{2.29}$$

This theoretical absorbed Dose ratio is also in good comparison to the DOS TEL measurements (2.32,2.35).

- An isotropic field could not explain the DOS TEL measurements during MTR 1.
- A comparison of DOS TEL data with results of a simple Monte-Carlo-Modell with different theta distribution showed:
- A $\cos^{2.5}$ theta distributed radiation field outside the ISS can explain the differences between single and telescope detectors for MTR DOS TEL during MTR1 mission phase.
- This anisotropy, changes the mean pathlength in the DOS TEL telescope from 1.21 to 1.17, which have a minor effect on the LET-spectra and so on mean Quality factors.

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Thank you for your attention.